





Chiral Nuclear Interactions for QMC Methods

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Neutron stars





https://physics.wustl.edu/quantum-monte-carlo-group



Quantum Monte Carlo Group for Nuclear Physics





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Microscopic calculations of nuclear systems

• OUR GOAL

Understand nuclear systems from a microscopic point of view, in terms of the interactions between individual nucleons and external probes

Nucleon-nucleon (NN) and 3N scattering data: "thousands" of experimental data available

Spectra, properties, and transition of nuclei: BE, radii, magnetic moments, beta decays rates, weak/ radiative captures, electroweak form factors, etc,...

Nucleonic matter equation of state: for ex. EOS neutron matter

• WHAT WE NEED!

Two and many-body interactions:

 $H = \sum_{i=1}^{A} \frac{\mathbf{p}_{i}^{2}}{2m_{i}} + \sum_{i< j=1}^{A} v_{ij} + \sum_{i< j< k=1}^{A} V_{ijk} + \dots$ $j^{\text{EW}} = \sum_{i=1}^{A} j_{ii} + \sum_{i< j=1}^{A} j_{ij} + \sum_{i< j< k=1}^{A} j_{ijk} + \dots$

i < i = 1 i < i < k = 1

VMC CVMC GFMC AFDMC (C)VMC $A \leq 12$ GFMC light systems 12AFDMC CVMC light to medium- $A \sim 50$ heavy nuclei AFDMC $E \rightarrow E_0$ AFDMC minimization infinite matter $A \to \infty$ τ propagation

Accurate many-body method:

Figure by Diego Lonardoni, LANL & MSU

Local chiral NN Hamiltonian with Δ's

 Local NN potentials including N2LO Δ-contributions and N3LO contacts have ben also devised and expressed as a sum of 16 spin-isospin operators



- Contact component parametrized by 26 LECs:
 - the functional form taken as $C_{R_S}(r) \propto e^{-(r/R_S)^2}$ with $R_S = 0.8~(0.7)~{
 m fm}$ a (b) models
 - models a (b) cutoff~500 MeV (600 MeV) in p-space

model	order	$E_{\rm Lab} ({\rm MeV})$	N_{pp+np}	χ^2/datum
Ia	N3LO	$0\!\!-\!\!125$	2668	1.05
Ib	N3LO	$0\!\!-\!\!125$	2665	1.07
IIa	N3LO	0–200	3698	1.37
IIb	N3LO	0 - 200	3695	1.37



MP et al. PRC 91, 024003 (2015); PRC 94, 054007 (2016)

Local chiral NNN Hamiltonian with **Δ's**

• Inclusion of 3N forces at N2LO: $V_{ijk} = V_{ijk}^{2\pi} + V_{ijk}^{CT}$

2) Fit to: $\triangleright E_0(^{3}\text{H}) = -8.482 \text{ MeV}$

C1 C3 C4

• GT m.e. in ³H β -decay

×--|

 $\underline{\mathbf{C}} \sim \tau_i \cdot \tau_j$

Model	c_D	c_E
Ia*	-0.635(255)	-0.09(8)
Ib^*	-4.705(285)	0.550(150)
IIa*	-0.610(280)	-0.350(100)
IIb*	-5.250(310)	0.05(180)



	Ia*	Ib*	IIa*	IIb*
c_D^*	-0.635	-4.71	-0.61	-5.25
c_E^*	-0.09	0.55	-0.35	0.05
LO	0.9272	0.9247	0.9261	0.9263
N2LO	0.0345	0.0517	0.0345	0.0515
N3LO(OPE)	0.0327	0.0454	0.0330	0.0465
N3LO(CT)	-0.0435	-0.0715	-0.0432	-0.0737

A.Baroni et al. PRC 98, 044003 (2018)

Ground state energies are ~5-6% away from exp data

- 1) Fit to: $\blacktriangleright E_0(^{3}\text{H}) = -8.482 \text{ MeV}$
 - $a_{nd} = (0.645 \pm 0.010) \text{ fm}$

Model	c_D	c_E
Ia	3.666	-1.638
Ib	-2.061	-0.982
IIa	1.278	-1.029
IIb	-4.480	-0.412



GT m.e. in ³H β -decay are ~3-4% away from exp data

Nuclear structure: spectra and radii of light-nuclei



1.4

- S.Gandolfi, D.Lonardoni, A.Lovato, MP Front. Phys. 8, 117 (2020)
 - Energy ratio between QMC results and experimental data (GFMC for NV2+3-Ia and AFDMC for GT+Eτ-1.0)
 - For NV2+3-Ia difference with experimental data less than 0.2 MeV/A, expected to be covered by truncation error estimate
 - For GT+Eτ-1.0 good agreement with experimental data within the theoretical uncertainty estimation.
 - Charge radii with respect to experimental data (GFMC for NV2+3-Ia and AFDMC for GT+Eτ-1.0)
 - Overall agreement with the experimental data for both models
 - For NV2+3-Ia, 9Li charge radius underpredicted, 12C slightly overestimated
 - For GT+Eτ-1.0, 6Li charge radius underpredicted (issue with AFDMC w.f.)

Nuclear structure: charge form factors

S.Gandolfi, D.Lonardoni, A.Lovato, MP Front. Phys. 8, 117 (2020)



Calculations using NV2+3-Ia are obtained with GFMC Calculations using GT+E τ -1.0 are obtained with AFDMC

Single-nucleon momentum distribution

• The probability of finding a nucleon with momentum **k** and spin-isospin projection in a given nuclear state

$$\rho_{\sigma\tau}(\mathbf{k}) = \int d\mathbf{r}_1' \, d\mathbf{r}_1 \, d\mathbf{r}_2 \cdots d\mathbf{r}_A \, \psi_{JM_J}^{\dagger}(\mathbf{r}_1', \mathbf{r}_2, \dots, \mathbf{r}_A) \, e^{-i\mathbf{k}\cdot(\mathbf{r}_1 - \mathbf{r}_1')} \, P_{\sigma\tau}(1) \, \psi_{JM_J}(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A)$$
The total normalization is: $N_{\sigma\tau} = \int \frac{d\mathbf{k}}{(2\pi)^3} \, \rho_{\sigma\tau}(\mathbf{k})$





Nucleon-pair momentum distribution

• The probability of finding two nucleons in a nucleus with relative momentum **q** and total-center-of-mass momentum **Q** in a spin-isospin projection

$$\rho_{ST}(\mathbf{q}, \mathbf{Q}) = \int d\mathbf{r}_1' d\mathbf{r}_1 d\mathbf{r}_2' d\mathbf{r}_2 d\mathbf{r}_3 \cdots d\mathbf{r}_A \psi_{JM_J}^{\dagger}(\mathbf{r}_1', \mathbf{r}_2', \mathbf{r}_3, \dots, \mathbf{r}_A) e^{-i\mathbf{q}\cdot(\mathbf{r}_{12} - \mathbf{r}_{12}')} \\ \times e^{-i\mathbf{Q}\cdot(\mathbf{R}_{12} - \mathbf{R}_{12}')} P_{ST}(12) \psi_{JM_J}(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \dots, \mathbf{r}_A)$$

The total normalization is:
$$N_{ST} = \int \frac{d\mathbf{q}}{(2\pi)^3} \frac{d\mathbf{Q}}{(2\pi)^3} \rho_{ST}(\mathbf{q}, \mathbf{Q})$$



Working on the connection with Short Time Approximation (STA) to account for two-body currents (S.Pastore et al.)

Multi-messenger Astronomy: EoS of Pure Neutron Matter



- Model dependence of the EOS at two-body level; AFDMC-UC calculations
- The max spread between AV18, NV2-IIa/IIb (fit to higher NN scattering data) is ~4 MeV per particle at $\rho=2\rho_0$
- Including NV2-Ia/Ib (fit to lower NN scattering data) the max spread is ~9MeV per particle at $\rho=2\rho_0$

MP et al. Phys. Rev. C 101, 045801 (2020) Editors' Suggestion

- Benchmark calculations between BHF, FHNC/ SOC, AFDMC(CP and UC)
- AFDMC-CP tends to overestimate the E/A compared to the AFDMC-UC: ~2-3MeV at $\rho=\rho_0$ and ~7-8MeV $\rho=2\rho_0$
- AFDMC-UC, BHF, FHNC/SOC are very close to each other up to $\rho=\rho_0\,({\rm \sim}1\,{\rm MeV})$
- FHNC/SOC is below AFDMC-UC and BHF at higher density; due to limited three-body terms into the cluster expansion



Model dependence very large in NV2+3 models: work in progress!

Optimization procedure for the LECs

- In EFT we inevitably end up with a model with parameters \mathbf{a}^{*} that we must fit to data

Least-square objective function for a set of observables

$$\chi^{2}(\mathbf{a}) = \sum_{i=1}^{N_{\text{data}}} \left(\frac{o_{i} - t_{i}(\mathbf{a})}{\delta o_{i}}\right)^{2}$$

"Conventional" least-square minimization:

 $\mathbf{a}^* = \min_{\mathbf{a}} \chi^2(\mathbf{a})$

- Take δo_i to be the experimental error (or same modification to take into account theoretical errors)
- Many optimization techniques suitable for this problem such as POUNDers, Newtons Methods,....
- UQ addressed as: Covariance methods, Bootstrapping, standard prescription truncation error, cutoff dependence,....
- over/under-fitting parameters,...

- O_i : measured values
- t_i : calculated values
- δo_i : uncertainty observables



- Particularly well suited for (any) EFT, but generally suited for theory errors
- Assumptions are made explicit (e.g. naturalness of LECs, truncation errors)
- Parameter estimation: conventional optimization recovered as special case
- Clear prescriptions for combining errors

MCMC Implementation and its application

- With MCMC, we can efficiently sample the parameter space to extract the posterior distribution
- A general MCMC works by:
 - 1. initialize walkers in the parameter space
 - 2. propose a new location for the walker to move to
 - 3.accept or reject the move based on the posterior of the current and proposed locations
 - 4. repeat 2. And 3. until the walkers converge to the final posterior
- In order to get familiarity with MCMC, we choose (for now) to work with a simpler case: only local short-range interactions

- To do so, we:
 - are using our existing codes written in Fortran to calculate the likelihood from NN scattering data
 - are using a pre-written MCMC package for the fitting: emcee package in Python
 - are using f2py to convert Fortan into a Python module



Results at NLO: Joint Probability Projections



J.Bub, MP, S.Pastore in progress

Conclusion/Outlook

- We are testing our models of NN+3N interactions with Δ-isobar based on chiral EFT framework in both light-nuclei and infinite nuclear matter
- We mainly focus our attention on studying static and dynamic properties of nuclei up to A=12 and EoS of infinite neutron matter
- We are working on constructing a coherent picture of lepton-nucleus interactions with particular focus on neutrino-nucleus scattering in a wide rage of energy and momenta
- With the current Delta-full chiral interactions/currents we are planning to investigate muon capture, neutrino scattering, neutrinoless double beta decay, beta-decays for A>10
- Inclusion of the present Delta-full 3N in calculations of infinite nucleonic matter; sensitivity studies of the EoS
- We are working on improvements of our nuclear models using Bayesian analysis tools to better access theoretical error estimation
- Particularly emphasis needs to be devoted to the 3N force; the formulation we have is too simplistic if we want to have a good descriptions of spectra, properties of light-nuclei, infinite nuclear matter, three-body observables with a certain degree of accuracy